

Differencing the Time-Dependent Schrödinger Equation

Now let's turn to the problem of solving the time-dependent Schrödinger equation:

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x)\psi.$$

Its form is basically that of the diffusion equation, except that the coefficients are complex and there is an additional “ $V\psi$ ” term on the right-hand side. The Crank-Nicolson scheme is still the method of choice. What's more, it can be shown that the Crank-Nicolson map from time n to time $n+1$ is *unitary*, so it preserves the normalization of ψ :

$$\int \psi^* \psi dx = 1.$$

As before, we derive the Crank-Nicolson differencing scheme by averaging together the fully explicit and fully implicit difference forms of the differential equation:

$$\begin{aligned} i\hbar \frac{\psi_j^{n+1} - \psi_j^n}{\Delta t} &= -\frac{\hbar^2}{2m} \frac{\psi_{j+1}^n - 2\psi_j^n + \psi_{j-1}^n}{\Delta x^2} + V_j \psi_j^n & (\text{explicit}), \\ i\hbar \frac{\psi_j^{n+1} - \psi_j^n}{\Delta t} &= -\frac{\hbar^2}{2m} \frac{\psi_{j+1}^{n+1} - 2\psi_j^{n+1} + \psi_{j-1}^{n+1}}{\Delta x^2} + V_j \psi_j^{n+1} & (\text{implicit}), \end{aligned}$$

where $V_j = V(x_j)$. Averaging and rearranging, we find

$$\begin{aligned} \psi_j^{n+1} &= \psi_j^n + \frac{1}{2}i\alpha \left\{ \psi_{j+1}^n - 2\psi_j^n + \psi_{j-1}^n + \psi_{j+1}^{n+1} - 2\psi_j^{n+1} + \psi_{j-1}^{n+1} \right\} \\ &\quad - \frac{1}{2}iV_j\Delta t(\psi_j^{n+1} + \psi_j^n)/\hbar, \end{aligned}$$

where $\alpha = \hbar\Delta t/(2m\Delta x^2)$. Moving the unknown “ $n+1$ ” terms to the left-hand side, we have

$$\begin{aligned} -\frac{1}{2}i\alpha \psi_{j-1}^{n+1} + \left[1 + i\alpha + \frac{1}{2}iV_j\Delta t/\hbar \right] \psi_j^{n+1} - \frac{1}{2}i\alpha \psi_{j+1}^{n+1} \\ = \frac{1}{2}i\alpha \psi_{j-1}^n + \left[1 - i\alpha - \frac{1}{2}iV_j\Delta t/\hbar \right] \psi_j^n + \frac{1}{2}i\alpha \psi_{j+1}^n. \end{aligned}$$

This complex matrix equation is then solved in exactly the same way as before.