

PHYS T580: The Standard Model

Homework #2

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1. **(SMIN 2.3)** Consider two particles of equal mass m connected by a spring of constant k and confined to move in one dimension. The entire system moves without friction. At equilibrium the spring has length L .

- (a) Write down the Lagrangian of this system as a function of x_1 and x_2 and their derivatives. Assume $x_2 > x_1$.

Solution:

The kinetic energy of each mass is $\frac{1}{2}m\dot{x}_1^2$ and for the second mass is $\frac{1}{2}m\dot{x}_2^2$. The total compression in the spring is $x_2 - x_1 - L$ so the total potential energy in the spring is $V = \frac{1}{2}k(x_2 - x_1 - L)^2$. So the Lagrangian of the system becomes

$$\mathcal{L} = \frac{1}{2}m\dot{x}_1^2 + \frac{1}{2}m\dot{x}_2^2 - \frac{1}{2}k(x_2 - x_1 - L)^2$$

This is the required Lagrangian. □

- (b) Write the Euler-Lagrange equation for this system.

Solution:

$$\begin{aligned} \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}_1} \right) &= \frac{\partial \mathcal{L}}{\partial x_1} &\Rightarrow m\ddot{x}_1 &= k(x_2 - x_1 - L) \\ \frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}_2} \right) &= \frac{\partial \mathcal{L}}{\partial x_2} &\Rightarrow m\ddot{x}_2 &= -k(x_2 - x_1 - L) \end{aligned}$$

These are the required Euler-Lagrange equations of the system. □

- (c) Make the change of variables

$$\Delta \equiv x_2 - x_1 - L \quad X = \frac{1}{2}(x_1 + x_2)$$

Write the Lagrangian in these new variables.

Solution:

Eliminating x_1 and x_2 between the two transformation variables Δ and X we get

$$\begin{aligned} x_2 - x_1 &= \Delta + L & x_1 + x_2 &= 2X \\ 2x_2 &= 2X + \Delta + L & \Rightarrow x_2 &= X + \frac{1}{2}(\Delta + L) & \dot{x}_2 &= \dot{X} + \frac{1}{2}\dot{\Delta} \\ 2x_1 &= 2X - \Delta - L & \Rightarrow x_1 &= X - \frac{1}{2}(\Delta + L) & \dot{x}_1 &= \dot{X} - \frac{1}{2}\dot{\Delta} \end{aligned}$$

Using these variables the lagrangian becomes

$$\begin{aligned}\mathcal{L} &= \frac{1}{2}m \left(\dot{X} + \frac{1}{2}\dot{\Delta} \right)^2 + \frac{1}{2}m \left(\dot{X} - \frac{1}{2}\dot{\Delta} \right)^2 + \frac{1}{2}k\Delta^2 \\ &= \frac{1}{2}m \left(2\dot{X}^2 + \dot{\Delta}^2 \right) + \frac{1}{2}k\Delta^2\end{aligned}$$

This is the lagrangian in the transformed coordinate system. \square

2. (SMIN 2.7) Show that the complex Lagrangian $\mathcal{L} = \partial_\mu \phi \partial^\mu \phi^* - m^2 \phi \phi^*$ is algebraically identical to

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi_1 \partial^\mu \phi_1 - \frac{1}{2} m_1^2 \phi_1^2 + \frac{1}{2} \partial_\mu \phi_2 \partial^\mu \phi_2 - \frac{1}{2} m_2^2 \phi_2^2$$

if $m_1 = m_2 = m$ and

$$\phi = \left(\frac{\phi_1 + \phi_2}{2} \right) + i \left(\frac{\phi_1 - \phi_2}{2} \right)$$

Solution:

Assuming the scalar fields ϕ_1 and ϕ_2 are real valued function. The complex field and its conjugate are

$$\begin{aligned}\phi &= \left(\frac{\phi_1 + \phi_2}{2} \right) + i \left(\frac{\phi_1 - \phi_2}{2} \right) & \phi^* &= \left(\frac{\phi_1 + \phi_2}{2} \right) - i \left(\frac{\phi_1 - \phi_2}{2} \right) \\ \Rightarrow \phi \phi^* &= \frac{1}{4} \left[(\phi_1 + \phi_2)^2 + (\phi_1 - \phi_2)^2 \right] = \frac{1}{4} (\phi_1^2 + \phi_2^2 + 2\phi_1\phi_2 + \phi_1^2 + \phi_2^2 - 2\phi_1\phi_2) \\ &= \frac{1}{2} (\phi_1^2 + \phi_2^2)\end{aligned}$$

$$\begin{aligned}\partial_\mu \phi \partial^\mu \phi^* &= \partial_\mu \left(\frac{\phi_1 + \phi_2}{2} + i \frac{\phi_1 - \phi_2}{2} \right) \partial^\mu \left(\frac{\phi_1 + \phi_2}{2} - i \frac{\phi_1 - \phi_2}{2} \right) \\ &= \frac{1}{2} (\partial_\mu \phi_1 + \partial_\mu \phi_2 + i \partial_\mu \phi_1 - i \partial_\mu \phi_2) \frac{1}{2} (\partial^\mu \phi_1 + \partial^\mu \phi_2 - i \partial^\mu \phi_1 + i \partial^\mu \phi_2) \\ &= \frac{1}{4} (\partial_\mu \phi_1 \partial^\mu \phi_1 + \partial_\mu \phi_1 \partial^\mu \phi_2 - i \partial_\mu \phi_1 \partial^\mu \phi_1 - i \partial_\mu \phi_1 \partial^\mu \phi_2 \\ &\quad + \partial_\mu \phi_2 \partial_\mu \phi_1 + \partial_\mu \phi_2 \partial^\mu \phi_2 - i \partial_\mu \phi_2 \partial^\mu \phi_1 - i \partial_\mu \phi_2 \partial^\mu \phi_2 \\ &\quad + i \partial_\mu \phi_1 \partial^\mu \phi_1 + i \partial_\mu \phi_1 \partial^\mu \phi_2 + \partial_\mu \phi_1 \partial^\mu \phi_1 + \partial_\mu \phi_1 \partial^\mu \phi_2 \\ &\quad + i \partial_\mu \phi_2 \partial^\mu \phi_1 + i \partial_\mu \phi_2 \partial^\mu \phi_2 + \partial_\mu \phi_2 \partial^\mu \phi_1 + \partial_\mu \phi_2 \partial^\mu \phi_2) \\ &= \frac{1}{2} (\partial_\mu \phi_1 \partial^\mu \phi_1 + \partial_\mu \phi_2 \partial^\mu \phi_2)\end{aligned}$$

Substituting these back in to the complex Lagrangian we get

$$\begin{aligned}\mathcal{L} &= \partial_\mu \phi \partial^\mu \phi^* - m^2 \phi \phi^* \\ &= \frac{1}{2} (\partial_\mu \phi_1 \partial^\mu \phi_1 + \partial_\mu \phi_2 \partial^\mu \phi_2) - m^2 \frac{1}{2} (\phi_1^2 + \phi_2^2) \\ &= \frac{1}{2} \partial_\mu \phi_1 \partial^\mu \phi_1 - \frac{1}{2} m_1^2 \phi_1^2 + \frac{1}{2} \partial_\mu \phi_2 \partial^\mu \phi_2 - \frac{1}{2} m_2^2 \phi_2^2\end{aligned}$$

This shows the two Lagrangian are equivalent. \square

3. (SMIN 2.9) Consider a lagrangian of real-valued scalar field:

$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2 - \frac{1}{6} c_3 \phi^3.$$

- (a) Is this Lagrangian Lorentz invariant? It is invariant under C , P , and T transformations individually?

Solution:

Since every term in the lagrangian is a scalar it is trivially Lorentz invariant. As ϕ is real valued scalar its complex conjugate is itself $\phi^* = \phi$ since the \hat{C} transformation transforms ϕ to ϕ^* which are identical so the Lagrangian is invariant under \hat{C} transformation.

It is not invariant under \hat{P} and \hat{T} transformation. □

- (b) What is the dimension of c_3 ?

Solution:

Since the dimension of Lagrangian density is $[E]^4$ and the dimensionality of ϕ is $[E]$ the dimensionality of c_3 is $[E]^3$ □

- (c) What is Euler-Lagrange equation for field?

Solution:

$$\begin{aligned}\partial_\mu \left(\frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi)} \right) &= \frac{\partial \mathcal{L}}{\partial \phi} \\ \partial_\mu (\partial^\mu \phi) &= -m^2 \phi - \frac{c_3}{2} \phi^2\end{aligned}$$

This is the required Euler-Lagrange equation for the given Lagrangian density. □

- (d) Ignoring the c_3 contribution, a free-field solution may be written

$$\phi_0(x) = Ae^{-ip \cdot x} + A^* e^{ip \cdot x}$$

for a complex coefficient A . Consider a lowest-order contribution for $\phi_1 \ll A$ to a perturbation such that $\phi(x) = \phi_0 + \phi_1$. Derive a dynamical equation for ϕ_1 .

Solution:

$$\partial_\mu \partial^\mu (\phi(x)) + m^2 \phi(x) + \frac{c_3}{2} \phi^2 = 0$$

Substituting $\phi = \phi_0 + \phi_1$

$$\begin{aligned}\Rightarrow \quad \partial_\mu \partial^\mu (\phi_0(x) + \phi_1(x)) + m^2 (\phi_0(x) + \phi_1(x)) + \frac{1}{2} c_3 (\phi_0 + \phi_1)^2 &= 0 \\ \Rightarrow \quad \partial_\mu \partial^\mu \phi_1(x) + m^2 \phi_1(x) + \frac{1}{2} c_3 \phi_0^2 \left(1 + \frac{\phi_1}{\phi_0} \right)^2 &= -\partial_\mu \partial^\mu \phi_0 - m^2 \phi_0 \\ \Rightarrow \quad \partial_\mu \partial^\mu \phi_1(x) + m^2 \phi_1(x) + \frac{1}{2} c_3 \phi_0^2 \left(1 + 2 \frac{\phi_1}{\phi_0} \right) &= -\partial_\mu \partial^\mu \phi_0 - m^2 \phi_0\end{aligned}$$

The first term in RHS of above expression is

$$\begin{aligned}\partial_\mu \partial^\mu \phi_0 &= g^{\mu\nu} \partial_\nu (\partial_\mu \phi_0) \\ &= g^{\mu\nu} \partial_\nu ((-ip_\mu) A e^{-ip \cdot x} + (ip_\mu) A^* e^{ip \cdot x}) \\ &= \partial_\nu ((-ip^\nu) A e^{-ip \cdot x} + (ip^\nu) A^* e^{ip \cdot x}) \quad (\text{Distributing } g^{\mu\nu}) \\ &= (-p^\nu p_\nu) A e^{-ip \cdot x} + (-p^\nu p_\nu) A^* e^{ip \cdot x} \\ &= -(E^2 - |\vec{p}|^2) \phi_0\end{aligned}$$

□