

PHYS 431: Galactic Astrophysics

Homework #4

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1. Estimate the masses of star clusters having

(a) root mean square velocity 10 km/s and half-mass radius 10 pc,

Solution:

Given $v_{rms} = 10 \text{ km/s}$, the mean square speed is $\langle v^2 \rangle = (10 \text{ km/s})^2 = 1 \times 10^4$. The total mass is given by

$$M = \frac{6R_h \langle v^2 \rangle}{G} = \frac{6 \cdot 1 \times 10^4 \cdot 10 \times 3.08 \times 10^{16}}{6.67 \times 10^{-11}} = 2.78 \times 10^{36} \text{ kg} = 1.39 \times 10^6 M_\odot$$

So the mass of the cluster is $1.39 \times 10^6 M_\odot$ □

(b) mean density 100 pc^{-3} , rms velocity 2 km/s , and mean stellar mass $0.8 M_\odot$,

Solution:

If the number density is n and average stellar mass is \bar{m} then the mean mass density

$$\rho = n \cdot \bar{m} = 100 \text{ pc}^{-3} \cdot 0.8 M_\odot = 80 M_\odot / \text{pc}^3; v_{rms} = 2 \text{ km/s} \Rightarrow \langle v^2 \rangle = 4 \times 10^4$$

. The density volume relation $\rho = \frac{3M}{4\pi R^3} \Rightarrow R = \left(\frac{3M}{4\pi\rho} \right)^{1/3}$.

$$M = \frac{6R \langle v^2 \rangle}{G} = \frac{6 \langle v^2 \rangle}{G} \left(\frac{3M}{4\pi\rho} \right)^{1/3} \Rightarrow M = \left(\frac{6 \langle v^2 \rangle}{G} \left(\frac{3}{4\pi\rho} \right)^{1/3} \right)^{3/2} = 4.53 \times 10^{34} \text{ kg} = 2.27 \times 10^4 M_\odot$$

□

(c) dynamical time $1 \times 10^6 \text{ yr}$ and radius 1 pc.

Solution:

The dynamical time $\tau = \left(\frac{3\pi}{G\rho} \right)^{1/2}$. Using $\rho = \frac{3M}{4\pi R^3}$ we get

$$M = \frac{4\pi^2 R^3}{G\tau^2} = 1.75 \times 10^{35} \text{ kg} = 8.79 \times 10^4 M_\odot$$

□

2. Interstellar gas in many galaxies is in virial equilibrium with the stars, in that the rms speed of the gas particles is the same as the rms stellar speed. Consider a large elliptical galaxy with a virial radius of 100 kpc and a mass of $1 \times 10^{12} M_\odot$ solar masses. Calculate the rms stellar velocity using the virial theorem. Hence estimate the temperature of the interstellar gas, assuming that it

is composed entirely of hydrogen.

Solution:

$$v_{rms} = \sqrt{\langle v^2 \rangle} = \left(\frac{GM}{6R} \right)^{\frac{1}{2}} = \left(\frac{6.67 \times 10^{-11} \cdot 1 \times 10^{12} \cdot 1.9 \times 10^{30}}{6 \cdot 1 \times 10^4 \cdot 3.08 \times 10^{16}} \right)^{\frac{1}{2}} = 2.68 \times 10^5 m/s = 268 km/s$$

The mass of hydrogen is $m_H = 1.67 \times 10^{-27} kg$. If all the interstellar mass was composed of hydrogen then the temperature would be given by reation

$$\frac{1}{2} m_H \langle v^2 \rangle = \frac{3}{2} kT \Rightarrow T = \frac{m_H \langle v^2 \rangle}{3k} = \frac{1.67 \times 10^{-27} \cdot (268 \times 10^3)^2}{3 \cdot 1.68 \times 10^{-23}} = 2.86 \times 10^6 K$$

□

3. Assuming an average stellar mass of $0.5M_\odot$ and $\Lambda = r_c/1AU$, lookup table values and find the relaxation time t_r at the center of globular cluster 47 Tucanae. Show that the crossing time $t_{cross} \approx 2r_c/\sigma_r \sim 1 \times 10^{-3} t_{relax}$

Solution:

The total number of stars in the cluster is given by

$$N = \frac{\text{Total Mass}}{\text{Mean Mass}} = \frac{800M_\odot}{0.5M_\odot} = 1600$$

The density of stars from table is $\rho = 10^{4.9} M_\odot/pc^3$. The dynamical time of the stars can be now calculated as

$$\tau = \left(\frac{GM}{r_c^3} \right)^{-\frac{1}{2}} = 3.09 \times 10^5 yr$$

.

Now the relaxation time

$$t_{relax} = \frac{N}{8.5 \ln(\Lambda)} \tau = \frac{1600}{8.5 \ln(r_c/1AU)} 3.09 \times 10^5 = 4.89 \times 10^6 yr$$

The cross time is

$$\frac{t_{cross}}{t_{relax}} = \frac{2r_c}{\sigma_r t_{relax}} = \frac{2 \cdot 0.7 pc}{1.1 \times 10^4 \cdot 4.89 \times 10^6} = 2.54 \times 10^{-2}$$

□

4. The velocities of stars in a stellar system are described by a three-dimensional Maxwellian distribution that is,

$$f(v) = Av^2 e^{-mv^2/2kT}$$

Here, A is a normalization constant, m is the stellar mass, assumed constant, k is Boltzmanns constant, and T is the temperature of the system. Verify the mean stellar kinetic energy is $\frac{1}{2} m \langle v^2 \rangle = \frac{3}{2} kT$

Solution:

The normalization condition gives

$$\int_0^{\infty} f(v)dv = \int_0^{\infty} Av^2 e^{-mv^2/2kT} dv = 1$$

To carry out the integration lets make some change of variables

$$\frac{mv^2}{2kT} = x; \Rightarrow v = \sqrt{\frac{2kT}{m}}x; \quad dv = \frac{kT}{mv}dx, \quad \text{As } v \rightarrow \{0, \infty\} \quad x \rightarrow \{0, \infty\}$$

Using these variable transformation, our normalization integral becomes.

$$A \int_0^{\infty} v^2 e^{-x} \frac{kT}{mv} dx = A \int_0^{\infty} \frac{kT}{m} \sqrt{\frac{2kT}{m}} x e^{-x} dx = A\sqrt{2} \left(\frac{kT}{m}\right)^{\frac{3}{2}} \int_0^{\infty} \sqrt{x} e^{-x} dx = 1$$

But by definition of gamma function $\Gamma(n) = \int_0^{\infty} x^{n-1} e^{-x}$ we get. And $\Gamma\left(\frac{3}{2}\right) = \frac{1}{2}\Gamma\left(\frac{1}{2}\right) = \frac{1}{2}\sqrt{\pi}$

$$A\sqrt{2} \left(\frac{kT}{m}\right)^{\frac{3}{2}} \int_0^{\infty} x^{\frac{3}{2}-1} e^{-x} dx = A\sqrt{2} \left(\frac{kT}{m}\right)^{\frac{3}{2}} \Gamma\left(\frac{3}{2}\right) = 1 \Rightarrow A = \frac{1}{\frac{1}{2}\sqrt{2\pi} \left(\frac{kT}{m}\right)^{\frac{3}{2}}}$$

The expectation value for the square of speed can be calculated as:

$$\langle v^2 \rangle = \int_0^{\infty} v^2 f(v)dv = A \int_0^{\infty} v^4 e^{-mv^2/2kT} dv$$

Carrying out same transformations as above we get.

$$\begin{aligned} \langle v^2 \rangle &= A \int_0^{\infty} 2^{\frac{3}{2}} \left(\frac{kT}{m}\right)^{\frac{5}{2}} x^{\frac{3}{2}} e^{-x} dx \\ &= A 2^{\frac{3}{2}} \left(\frac{kT}{m}\right)^{\frac{5}{2}} \int_0^{\infty} x^{\frac{5}{2}-1} e^{-x} dx = A 2^{\frac{3}{2}} \left(\frac{kT}{m}\right)^{\frac{5}{2}} \Gamma\left(\frac{5}{2}\right) = A 2^{\frac{3}{2}} \left(\frac{kT}{m}\right)^{\frac{5}{2}} \frac{3}{4} \sqrt{\pi} \\ &= \frac{1}{\frac{1}{2}\sqrt{2\pi} \left(\frac{kT}{m}\right)^{\frac{3}{2}}} \times 2^{\frac{3}{2}} \left(\frac{kT}{m}\right)^{\frac{5}{2}} \frac{3}{4} \sqrt{\pi} \\ &= \frac{3kT}{m} \\ \Rightarrow \frac{1}{2}m\langle v^2 \rangle &= \frac{3}{2}kT \end{aligned}$$

So the kinetic energy of each mass is $\frac{3}{2}kT$ if the velocity distribution of the ensemble of mass follow Maxwellian distribution function. \square

5. Work out the details of the simple evaporative model discussed in class. Stars evaporate from a cluster of mass M on a time scale $t_{ev} = \alpha t_R$, where $\alpha \gg 1$, so

$$\frac{dM}{dt} = -\frac{M}{\alpha t_R} \quad (1)$$

For pure evaporation, each escaping star carries off exactly zero energy (i.e. stars barely escape the cluster potential), so the total energy of the cluster remains constant.

- (a) If the cluster potential energy can always be written as $U = -k\frac{GM^2}{2R}$ for fixed k , where R is a characteristic cluster radius, and assuming that the cluster is always in virial equilibrium, show that $R \propto M^2$ as the cluster evolves.

Solution:

The potential energy relation can be reorganized as

$$R = -\frac{kG}{2U}M^2; \quad \Rightarrow R = \beta_0 M^2; \quad \Rightarrow R \propto M^2; \quad \text{Where } \beta_0 = -\frac{kG}{2U}$$

So $R \propto M^2$. □

- (b) Assuming that the relaxation time t_R scales as $M^{1/2}R^{3/2}$ so

$$t_R = t_{R0} \left(\frac{M}{M_0}\right)^{1/2} \left(\frac{R}{R_0}\right)^{3/2} \quad (2)$$

Solve (??) to determine the lifetime of the cluster (in terms of its initial relaxation time t_{R0}). Also write down an expression for the mean cluster density as a function of time.

Solution:

We can write Eq. (??) as $t_R = \beta_1 M^{1/2} R^{3/2}$. Since $R = \beta_0 M^2$. We now have, $t_R = \beta_1 M^{1/2} (\beta_0 M^2)^{3/2}$;

$$\Rightarrow t_R = \beta_3 M^{7/2}$$

Suppose T is the lifetime of the cluster that had initial mass of M_i then as time goes from 0 to T mass goes from M_i to 0. Using t_R in Eq. (??) we get

$$\frac{dM}{dt} = -\frac{1}{\alpha} \frac{M}{\beta_3 M^{7/2}}; \Rightarrow \int_{M_i}^0 M^{5/2} dM = -\beta_4 \int_0^T dt; \Rightarrow -\frac{2}{7} M_i^{7/2} = -\beta_4(T) \Rightarrow T \propto M_i^{7/2}$$

So the lifetime of the cluster is $T \propto M_i^{7/2}$ ■.

Now the density $\rho \propto \frac{M}{R^3}$. But for a system in dynamical equilibrium we have $R \propto M^2$. This gives $\rho \propto \frac{M}{(M^2)^3} = M^{-5} \Rightarrow M \propto \rho^{-5}$ Eq. (??) can be solved as a function of time as above and written as

$$M = \beta_5 t^{2/7} \Rightarrow M^{-5} = \beta_5 t^{-10/7} \Rightarrow \rho = \beta_6 t^{-10/7}$$

□

- (c) Estimate this for a globular cluster of mass $5 \times 10^5 M_\odot$ radius $10pc$ and mean stellar mass $0.5M_\odot$

Solution:

The density of this cluster is

$$\rho \approx \frac{M}{R^3} = \frac{5 \times 10^5 M_\odot}{10^3 pc^3} = 9.86 \times 10^{-14} kg/m^3 = 5 \times 10^2 M_\odot/pc^3$$

. The number of star is

$$N = \frac{M_{tot}}{m_{av}} = \frac{5 \times 10^5 M_\odot}{0.5 M_\odot} = 1 \times 10^6$$

The time scale then is

$$t = \left(\frac{GM}{R^3} \right)^{-1/2} = \frac{6.67 \times 10^{-11} \cdot 5 \times 10^5 M_\odot}{10^3 pc^3} = 6.67 \times 10^5 yr$$

□